

# All-optical modulation in a CMOS-compatible amorphous silicon-based device

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Active silicon photonic devices, which dynamically control the flow of light, have received significant attention for their use in on-chip optical networks. High-speed active silicon photonic modulators and switches rely on the plasma dispersion effect, where a change in carrier concentration causes a variation in the refractive index. The necessary electron and hole concentration change can be introduced either by optical pumping, or by direct electrical injection and depletion. We demonstrate a fast photoinduced absorption effect in low loss hydrogenated amorphous silicon (a-Si:H) waveguides deposited at a temperature as low as 190 °C. Significant modulation ( $M_{\%}$ , ~90%) occurs with a 1 mm-long device. We attribute the enhanced modulation to the significantly larger free-carrier absorption effect of a-Si:H. The complementary metal-oxide semiconductor (CMOS) compatible technology of a-Si:H could be considered as a promising candidate to enable an easy back-end integration with standard microelectronics processes.

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## 1 INTRODUCTION

Photonics is a rapidly growing sector in the global economy. Optical communications, optical storage, imaging, lighting, optical sensors or security are just a few examples. Even if photonics could bring new functionalities to electronic components, *e.g.* low propagation losses, high bandwidth, wavelength multiplexing and immunity to electromagnetic noise, the high cost of photonic components is a major obstacle to their deployment in most of the application fields. In microelectronics, many applications can be realised in a much more compact and cost-effective way by integrating the required functionality in a single chip. Photonics on CMOS seems to be the way to tackle such issue by developing a small number of generic integration technologies with a level of functionality that can address a broad range of applications [1].

Hydrogenated amorphous silicon (a-Si:H), deposited using the CMOS-compatible low temperature (120–400 °C) Plasma Enhanced Chemical Vapour Deposition (PECVD) technique, is recently emerging as a useful material for realising minimally-invasive on-chip passive [2]–[4] and active [5]–[8] devices for optical interconnects. In this paper we report experimental results on all-optical modulation in a compact device. Modulation is achieved by using short optical pump

laser pulses to induce an absorption change in the a-Si:H waveguide.

## 2 WAVEGUIDE STRUCTURE AND FABRICATION

A schematic cross section of the designed waveguide is shown in Figure 1. The device, realised on a p-doped crystalline silicon (c-Si) substrate ( $\rho=0.001 \Omega \cdot \text{cm}$ ), consists of a rib waveguide made of a 3- $\mu\text{m}$ -thick a-Si:H undoped layer (refractive index  $n_{a\text{-Si:H}}=3.58$ ) on a 400 nm-thick  $\text{SiO}_2$  low refractive index cladding ( $n_{\text{SiO}_2}=1.45$ ) ensuring a vertical optical confinement. The undoped a-Si:H layer is deposited from the plasma-assisted decomposition of  $\text{SiH}_4$  at a flow of 20 sccm (standard cubic centimetres per minute) and an RF power of  $P_{RF}=4 \text{ W}$  at the frequency of  $f=13.56 \text{ MHz}$ . The maximum temperature during the process was as low as 190 °C, which makes it fully compatible with standard CMOS processes. The CMOS-friendly technological deposition process allows, in principle, the simple fabrication of a photonic layer on top of the electronic microchip at the end of the CMOS flow, just before encapsulation.

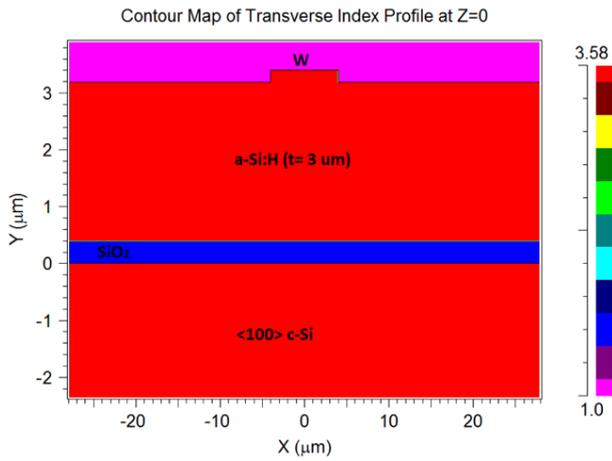


FIG. 1 Schematic cross section of the realised waveguides and corresponding refractive index colour bar. The crystalline silicon substrate is 300 μm thick.

Standard technological processes have been used, *i.e.* optical lithography and reactive ion etching (RIE), to realise the rib waveguide. Optical simulations [9] show that the designed waveguide supports only one mode for TE and TM polarisation once a 8 μm-wide and 210 nm-high rib is defined while the first-order higher modes, TE<sub>01</sub> and TM<sub>01</sub>, are localised in the slabs, far from the rib.

### 3 ALL-OPTICAL EXPERIMENTAL RESULTS

We have tested the behavior of the a-Si:H-based waveguiding active device by inducing a strong photoinduced absorption on a 1.0 mm-long device using a pump-probe optical setup. In such configuration, while the pump beam is orthogonal to the film, the probe beam is launched transversally to the film thickness and guided by the waveguide core, increasing in this way the interaction length between the active material and the optical signal.

In our experiments, the waveguide propagation losses were first measured by the cut-back technique. Several samples with lengths in the range from 800 to 1500 μm were cut from the substrate by cleavage. Sample facets did not receive polishing treatment. From these samples we calculated a loss coefficient of 1.3±0.1 dB/cm (0.29±0.02 cm<sup>-1</sup>). The overall insertion losses of the devices are ~10 dB with the main contribution due to the coupling losses, a drawback of the high refractive index of a-Si:H used as core material.

The experimental setup is depicted in Figure 2. A 4 mW laser diode probe radiation at the wavelength of 1550 nm was coupled into the device via a lensed fiber. The transmitted light was collected at the chip output by a multi-mode fiber and detected by a 2 GHz bandwidth InGaAs photodiode. We used the 2<sup>nd</sup> harmonic (λ<sub>PUMP</sub>=532 nm) of a high energy Nd:YAG laser for inducing the free carrier injection in the waveguide. The low level optical probe signal at the waveguide output was recovered by mean of an erbium doped fiber amplifier (EDFA).

Optical modulation is achieved by using short optical pump

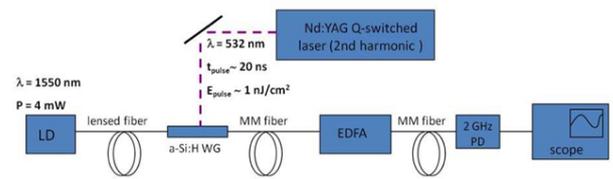


FIG. 2 Experimental setup used for the characterisation of the photoinduced absorption effect in an a-Si:H layer.

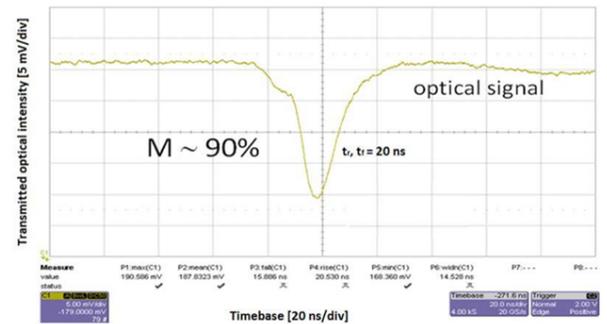


FIG. 3 Photoinduced absorption effect induced in a 3 μm-thick, 1-mm-long, a-Si:H rib waveguide. The electrical signal at the photodiode output is shown in the figure. The pump pulse was a high-energy, 532 nm wavelength, 20-ns-long, coherent light pulses. The measured rise and fall times are ~20 ns and the modulation depth is 90%.

pulses ( $E_{pulse} \sim 1 \text{ nJ/cm}^2$ ) to induce a change in the extinction coefficient,  $k$ , of the amorphous waveguiding material. We attributed the continuous-wave laser modulation to the free carriers absorption effect induced by the generated photo-excited carriers (the quasi-bandgap of a-Si:H is ~1.75 eV). We measured a modulation depth, defined as:  $(I_{MAX}-I_{MIN})/I_{MAX}$  where  $I_{MAX}$  and  $I_{MIN}$  are the maximum and minimum intensities of the transmitted signal, of  $M_{\%} \sim 90\%$  when an incident pump pulse is delivered. Due to the poor overlap of the pump light with the a-Si:H active region, which is mainly confined under the rib waveguide (Active Area~8 μm×1 mm) where the probe beam propagates, we estimated that only 80 fJ of the pump energy contributes to the absorption effect. Consequently, we are confident that the modulation depth could be improved with a better overlap even providing a lower pump pulse energy. It is worth noting that no Fabry-Perot interference effect from the ends facets were observed in transmission measurements.

As shown in Figure 3, the probe optical beam at the photodiode output falls when the pump photons are absorbed. In fact, since the photon energy is much larger than the quasi-bandgap energy of a-Si:H, the electrons are excited into the conduction band highly increasing the free carrier absorption effect. We measured rise and fall time of the order of 20 ns, which are higher than what observed in c-Si, where characteristic times of the order of hundreds of picoseconds have been observed [10]. We note, however, that in our experiments the characteristic times are probably limited by the dynamics of our Nd:YAG pump beam. From the degree of modulation measured in our experiment, we estimated an absorption coefficient variation  $\Delta\alpha \sim 23 \text{ cm}^{-1}$ .

The refractive index change in silicon, due to injection or depletion of free carrier at λ=1.55 μm, can be derived to a first

order approximation from the classical Drude model also for amorphous silicon semiconductor [11], in particular the absorption coefficient variation is given by:

$$\Delta\alpha = \frac{e^3\lambda^2}{4\pi^2c^3\varepsilon_0n} \left( \frac{\Delta N_e}{m_e^2\mu_e} + \frac{\Delta N_h}{m_h^2\mu_h} \right) \quad (1)$$

In Eq. (1),  $e$  is the electron charge,  $\lambda$  is the probe wavelength,  $\varepsilon_0$  is permittivity of free space,  $n_0$  is the refractive index of the material,  $m_e$  and  $m_h$  are the effective masses of electrons and holes,  $\mu_e$  and  $\mu_h$  are the mobilities of the carriers in a-Si:H. Substituting for  $m_e=0.5\cdot m_0$ ,  $m_h=1.0\cdot m_0$ ,  $m_0=9.1\cdot 10^{-31}$  kg,  $\mu_e=2.0$  cm<sup>2</sup>/V·s,  $\mu_h=0.4$  cm<sup>2</sup>/V·s [11] in Eq. (1), yields a theoretically estimated average free carriers concentration change of:

$$\Delta N_e \sim \Delta N_h = 1.4 \times 10^{17} \text{ cm}^{-3}.$$

## 4 CONCLUSION

We have demonstrated an all-optical absorption-induced modulation of light in an amorphous silicon waveguide based on the free carrier effect using a pump-probe technique. The optical signal turn-on/turn off transients were measured to be of  $\sim 20$  ns, limited however by the pump laser source dynamic characteristics. The modulation depth is 90% when an incident pump pulse with energy of  $E_{pump} \sim 80$  fJ is absorbed in the active waveguiding region.

Technologically, the fabrication process involves temperatures below 190 °C, which are compatible with standard microelectronic processes and in particular suitable for the realisation of a photonic layer on top of an integrated circuit.

## 5 ACKNOWLEDGEMENTS

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